

I'm not a robot



Force Centrifugal force reactive Coriolis force Pendulum Tangential speed Tangential speed Tangential speed frequency acceleration / displacement / velocity Scientists Kepler Huygens Newton Horrocks Halley Maupertuis Daniel Bernoulli Johann Bernoulli Euler d'Alembert Clairaut Laplace Poisson Hamilton Jacobi Cauchy Routh Liouville Appellet Gibbs Koopman van Neumann Physics portal Categrory: In mechanics and physics, simple harmonic motion (sometimes abbreviated as SHM) is a special type of periodic motion an object experiences by means of a restoring force whose magnitude is directly proportional to the distance of the object from an equilibrium position and acts towards the equilibrium position. It results in an oscillation that is described by a sinusoid which continues indefinitely (if uninhibited by friction or any other dissipation of energy).[1] Simple harmonic motion can serve as a mathematical model for a variety of motions, but is typified by the oscillation of a mass on a spring when it is subject to the linear elastic restoring force given by Hooke's law. The motion is sinusoidal in time and demonstrates a single resonant frequency. Other phenomena can be modeled by simple harmonic motion, including the motion of a simple pendulum, although for it to be an accurate model, the net force on the object at the end of the pendulum must be proportional to the displacement (and even so, it is only a good approximation when the angle of the swing is small; see small-angle approximation). Simple harmonic motion can also be used to model molecular vibration. A mass-spring system is a classic example of simple harmonic motion. Simple harmonic motion provides a basis for the characterization of more complicated periodic motion through the techniques of Fourier analysis. The motion of a particle moving along a straight line with an acceleration whose direction is always toward a fixed point on the line and whose magnitude is proportional to the displacement from the fixed point is called simple harmonic motion.[2] In the diagram, a simple harmonic oscillator, consisting of a weight attached to one end of a spring, is shown. The other end of the spring is connected to a rigid support such as a wall. If the system is left at rest at the equilibrium position then there is no net force acting on the mass. However, if the mass is displaced from the equilibrium position, the spring exerts a restoring elastic force that obeys Hooke's law. Mathematically,

F
=
−
k
x
,

{\displaystyle \mathbf {F} =-k\mathbf {x} ,}

 where **F** is the restoring elastic force exerted by the spring (in SI units: N), *k* is the spring constant (N·m−1), and *x* is the displacement from the equilibrium position (in metres). For any simple mechanical harmonic oscillator: When the system is displaced from its equilibrium position, a restoring force that obeys Hooke's law tends to restore the system to equilibrium. Once the mass is displaced from its equilibrium position, it experiences a net restoring force. As a result, it accelerates and starts going back to the equilibrium position. When the mass moves closer to the equilibrium position, the restoring force decreases. At the equilibrium position, the net restoring force vanishes. However, at *x* = 0, the mass has momentum because of the acceleration that the restoring force has imparted. Therefore, the mass continues past the equilibrium position, compressing the spring. A net restoring force then slows it down until its velocity reaches zero, whereupon it is accelerated back to the equilibrium position again. As long as the system has no energy loss, the mass continues to oscillate. Thus simple harmonic motion is a type of periodic motion. If energy is lost in the system, then the mass exhibits damped oscillation. Note if the real space and phase space plot are not co-linear, the phase space motion becomes elliptical. The area enclosed depends on the amplitude and the maximum momentum. In Newtonian mechanics, for one-dimensional simple harmonic motion, the equation of motion, which is a second-order linear ordinary differential equation with constant coefficients, can be obtained by means of Newton's second law and Hooke's law for a mass on a spring.

F

n
e
t

=
m

d

2

x

d

t

2

=
−
k
x
,

{\displaystyle {\mathrm {net} }=m{\frac {\mathrm {d} ^{2}x}{\mathrm {d} t^{2}}}}=-kx,}

 where *m* is the inertial mass of the oscillating body, *x* is its displacement from the equilibrium (or mean) position, and *k* is a constant (the spring constant for a mass on a spring). Therefore,

d

2

x

d

t

2

=
−
k
m
x

{\displaystyle {\frac {\mathrm {d} ^{2}x}{\mathrm {d} t^{2}}}}=-{\frac {k}{m}}x}

 Solving the differential equation above produces a solution that is a sinusoidal function:

x
(
t
)
=

c

1

cos
⁡
(
ω
t
)
+

c

2

sin
⁡
(
ω
t
)
,

{\displaystyle x(t)=c_{1}\cos {\left(\omega t\right)}+c_{2}\sin {\left(\omega t\right)},}

 where

ω
=
k

/

m
.

{\textstyle {\omega }={\sqrt {\tfrac {k}{m}}}.}

 The meaning of the constants

c

1

 and

c

2

(

{\displaystyle c_{1}}

 and

c

2

(

{\displaystyle c_{2}}

 can be easily found: setting

t
=
0

{\displaystyle t=0}

 on the equation above we see that

x
(
0
)
=

c

1

{\displaystyle x(0)=c_{1}}

, so that

c

1

(

{\displaystyle c_{1}}

 is the initial position of the particle,

c

1

=
x
(
0
)

{\displaystyle c_{1}=x_{0}}

; taking the derivative of that equation and evaluating at zero we get that

x
˙

(
0
)
=

ω

c

2

{\displaystyle {\dot {x}}(0)=\omega c_{2}}

, so that

c

2

(

{\displaystyle c_{2}}

 is the initial speed of that particle divided by the angular frequency,

c

2

=

v

0

ω

{\displaystyle c_{2}={\frac {v_{0}}{\omega }}}

. Thus we can write:

x
(
t
)
=

x

0

cos
⁡
(
k
m
t
)
+

v

0

k
m

sin
⁡
(
k
m
t
)
,

{\displaystyle x(t)=x_{0}\cos {\sqrt {\frac {k}{m}}}t\right)+{\frac {v_{0}}{k}}\sin {\sqrt {\frac {k}{m}}}t\right).}

 This equation can also be written in the form:

x
(
t
)
=
A
cos
⁡
(
ω
t
−
ϕ
)
,

{\displaystyle x(t)=A\cos {\left(\omega t-\varphi \right)},}

 where

A
=

c

1

2

+

c

2

2

{\displaystyle A={\sqrt {{c_{1}}^{2}+{c_{2}}^{2}}}}

 tan

ϕ
=

c

2

c

1

{\displaystyle \tan \varphi ={\frac {c_{2}}{c_{1}}}}

;

sin
⁡
ϕ
=

c

2

A
,
cos
⁡
ϕ
=

c

1

A

{\displaystyle \sin \varphi ={\frac {c_{2}}{A}},\;\cos \varphi ={\frac {c_{1}}{A}}}

 or equivalently

A
=

|

c

1

+

c

2

i

|

,

{\displaystyle A=|c_{1}+c_{2}i|,\;\varphi =\arg {\,c_{1}+c_{2}i}}

(

{\displaystyle \varphi =\arg(c_{1}+c_{2}i)}

 In the solution,

c

1

 and

c

2

 are two constants determined by the initial conditions (specifically, the initial position at time

t
=
0

{\displaystyle t=0}

 is

c

1

, while the initial velocity is

c

2

ω

{\displaystyle c_{2}\omega }

), and the origin is set to be the equilibrium position.[A] Each of these constants carries a physical meaning of the motion: *A* is the amplitude (maximum displacement from the equilibrium position),

ω
=
2
π
f

{\displaystyle \omega =2\pi f}

 is the angular frequency, and

ϕ
 is the initial phase.[B] Using the techniques of calculus, the velocity and acceleration as a function of time can be found:

v
(
t
)
=

d

x

d

t

=
−
A
ω
sin
⁡
(
ω
t
−
ϕ
)
,

{\displaystyle a(t)={\frac {\mathrm {d} ^{2}x}{\mathrm {d} t^{2}}}=-A\omega ^{2}\cos(\omega t-\varphi).}

 Maximum acceleration:

a

ω

2

tan
⁡
(
at
extreme
points
)

{\displaystyle \tan \varphi ={\frac {c_{1}}{c_{2}}}}

 Since

ω
=
2
π
f
,
f
=

1

2
π
k
m

,

{\displaystyle \omega ^{2}={\frac {k}{m}}}

 Since

ω
=
2
π
f
,
f
=

1

2
π
k
m

,

{\displaystyle \omega ^{2}={\frac {k}{m}}}

 and, since

T
=

1

f

 where *T* is the time period,

T
=

2
π
m
k

{\displaystyle T=2\pi {\sqrt {\frac {m}{k}}}}

. These equations demonstrate that the simple harmonic motion is isochronous (the period and frequency are independent of the amplitude and the initial phase of the motion). Substituting

ω
2

 with

k
m

, the kinetic energy *K* of the system at time *t* is

K
(
t
)
=

1

2

m

v

2

(
t
)
=

1

2

m

ω

2

x

2

sin
⁡
(
ω
t
−
ϕ
)
=

1

2

k

A

2

sin
⁡
(
ω
t
−
ϕ
)
,

{\displaystyle K(t)={\frac {1}{2}}m{\dot {x}}^{2}(t)={\frac {1}{2}}m\omega ^{2}A^{2}\sin ^{2}(\omega t-\varphi)={\frac {1}{2}}kA^{2}\sin ^{2}(\omega t-\varphi),}

 and the potential energy is

U
(
t
)
=

1

2

k

x

2

(
t
)
=

1

2

k

A

2

cos
⁡
(
ω
t
−
ϕ
)
,

{\displaystyle U(t)={\frac {1}{2}}kx^{2}(t)={\frac {1}{2}}kA^{2}\cos ^{2}(\omega t-\varphi).}

 In the absence of friction and other energy loss, the total mechanical energy has a constant value

E
=
K
+
U
=

1

2

k

A

2

.

{\displaystyle E=K+U={\frac {1}{2}}kA^{2}.}

 An undamped spring-mass system undergoes simple harmonic motion. The following physical systems are some examples of simple harmonic oscillator. A mass *m* attached to a spring of spring constant *k* exhibits simple harmonic motion in closed space. The equation for describing the period:

T
=

2
π
m
k

{\displaystyle T=2\pi {\sqrt {\frac {m}{k}}}}

 shows the period of oscillation is independent of the amplitude, though in practice the amplitude should be small. The above equation is also valid in the case when an additional constant force is being applied on the mass, i.e. the additional constant force cannot change the period of oscillation. Simple harmonic motion can be considered the one-dimensional projection of uniform circular motion. If an object moves with angular speed

ω
 around a circle of radius *r* centered at the origin of the xy-plane, then its motion along each coordinate is simple harmonic motion with amplitude *r* and angular frequency

ω
. The motion of a body in which it moves to and from a definite point is also called oscillatory motion or vibratory motion. The time period is able to be calculated by

T
=

2
π
l
g

{\displaystyle T=2\pi {\sqrt {\frac {l}{g}}}}

 where *l* is the distance from rotation to the object's center of mass undergoing SHM and *g* is gravitational acceleration. This is analogous to the mass-spring system. A pendulum making 25 complete oscillations in 60 s, a frequency of 0.416 Hertz. In the small-angle approximation, the motion of a simple pendulum is approximated by simple harmonic motion. The period of a mass attached to a pendulum of length *l* with gravitational acceleration *g*

(

{\displaystyle g}

 is given by

T
=

2
π
l
g

{\displaystyle T=2\pi {\sqrt {\frac {l}{g}}}}

. This shows that the period of oscillation is independent of the amplitude and mass of the pendulum but not of the acceleration due to gravity,

g

{\displaystyle g}

, therefore a pendulum of the same length on the Moon would swing more slowly due to the Moon's lower gravitational field strength. Because the value of

g

{\displaystyle g}

 varies slightly over the surface of the earth, the time period will vary slightly from place to place and will also vary with height above sea level. This approximation is accurate only for small angles because of the expression for angular acceleration *a* being proportional to the sine of the displacement angle:

−
m
g
l
sin
⁡
θ
=
l
α
,

{\displaystyle -mg\sin \theta =l\alpha ,}

 where *l* is the moment of inertia. When

θ
 is small,

sin
⁡
θ
=
θ
 and therefore the expression becomes

−
m
g
l
θ
=
l
α

{\displaystyle -mg\theta =l\alpha }

 which makes angular acceleration directly proportional and opposite to

θ
, satisfying the definition of simple harmonic motion (that net force is directly proportional to the displacement from the mean position and is directed towards the mean position). Main article: Scotch yoke A Scotch yoke mechanism can be used to convert between rotational motion and linear reciprocating motion. The linear motion can take various forms depending on the shape of the slot, but the basic yoke with a constant rotation speed produces a linear motion that is simple harmonic in form. Scotch yoke animation Circle group Complex harmonic motion Damping ratio Harmonic oscillator Isochronous timing Lorentz oscillator model Newtonian mechanics Pendulum Rayleigh-Lorentz pendulum Small-angle approximation String vibration Uniform circular motion

^
 The choice of using a cosine in this equation is a convention. Other valid formulations are:

x
(
t
)
=
A
sin
⁡
(
ω
t
+
ϕ
′
)
,

{\displaystyle x(t)=A\sin {\left(\omega t+\varphi 'right)},}

 where

tan
⁡
ϕ
′
=

c

1

c

2

,

{\displaystyle \tan \varphi '={\frac {c_{1}}{c_{2}}}}

 since

cos
⁡
θ
=
sin
⁡
(
2
π
−
θ
)
,

{\displaystyle \cos \theta =\sin(2\pi -\theta),}

 and thus when

x

max

=
A
,

{\displaystyle x_{\mathrm {max} }=A,}

 Simple harmonic motion [Formula, Examples, & Facts | Britannica". britannica.com. 2024-09-30. Retrieved 2024-10-11. "Simple Harmonic Motion - Concepts". Fowles, Grant R., Cassiday, George L. (2005). Analytical Mechanics (7th ed.). Thomson Brooks/Cole. ISBN 0-534-49492-7. Taylor, John R. (2005). Classical Mechanics. University Science Books. ISBN 1-891389-22-X. Thornton, Stephen T.; Marion, Jerry B. (2003). Classical Dynamics of Particles and Systems (5th ed.). Brooks Cole. ISBN 0-534-40896-6. Walker, Jearl (2011). Principles of Physics (9th ed.). Hoboken, New Jersey: Wiley. ISBN 978-0-470-56158-4. Wikimedia Commons has media related to simple harmonic motion. Simple Harmonic Motion from HyperPhysics Java simulation of spring-mass oscillator Geogebra applet for spring-mass, with 3 attached PDFs on SHM, driven/damped oscillators, spring-mass with friction Retrieved on " Download the Testbook APP & Get Pass Pro Max FREE for 7 Days10.00+ Study NotesRealtime Doubt Support71000+ Mock TestsRankers Test Series+ more benefitsDownload App Now By the end of this section, you will be able to: Describe a simple harmonic oscillator. Explain the link between simple harmonic motion and waves. The oscillations of a system in which the net force can be described by Hooke's law are of special importance, because they are very common. They are also the simplest oscillatory systems. Simple Harmonic Motion (SHM) is the name given to oscillatory motion for a system where the net force can be described by Hooke's law, and such a system is called a simple harmonic oscillator. If the net force can be described by Hooke's law and there is no damping (by friction or other non-conservative forces), then a simple harmonic oscillator will oscillate with equal displacement on either side of the equilibrium position, as shown for an object on a spring in Figure 1. The maximum displacement from equilibrium is called the amplitude *X*. The units for amplitude and displacement are the same, but depend on the type of oscillation. For the object on the spring, the units of amplitude and displacement are meters; whereas for sound oscillations, they have units of pressure (and other types of oscillations have yet other units). Because amplitude is the maximum displacement, it is related to the energy in the oscillation. Figure 1. An object attached to a spring sliding on a frictionless surface is an uncomplicated simple harmonic oscillator. When placed from equilibrium, the object performs simple harmonic motion—that has an amplitude *X* and a period *T*. The object's maximum speed occurs as it passes through equilibrium. The stiffer the spring is, the smaller the period *T*. The greater the mass of the object is, the greater the period *T*. Find a bowl or basin that is shaped like a hemisphere on the inside. Place a marble inside the bowl and tilt the bowl periodically so the marble rolls from the bottom of the bowl to equally high points on the sides of the bowl. Get a feel for the force required to maintain this periodic motion. What is the restoring force and what role does the force you apply play in the simple harmonic motion (SHM) of the marble? What is so significant about simple harmonic motion? One special thing is that the period *T* and frequency *f* of a simple harmonic oscillator are independent of amplitude. The string of a guitar, for example, will oscillate with the same frequency whether plucked gently or hard. Because the period is constant, a simple harmonic oscillator can be used as a clock. Two important factors do affect the period of a simple harmonic oscillator. The period is related to how stiff the system is. A very stiff object has a large force constant *k*, which causes the system to have a smaller period. For example, you can adjust a diving board's stiffness—the stiffer it is, the faster it vibrates, and the shorter its period. Period also depends on the mass of the oscillating system. The more massive the system is, the longer the period. For example, a heavy person on a diving board bounces up and down more slowly than a light one. In fact, the mass *m* and the force constant *k* are the only factors that affect the period and frequency of simple harmonic motion. The period of a simple harmonic oscillator is given by

[
a
t
e
x
]
T
=

2
π

s
q
r
t

(

f
r
a
c
m
k

)

{\displaystyle T=2\pi {\sqrt {\frac {m}{k}}}}\!\!

 and, because

[
a
t
e
x
]
f
=

1
T

{\displaystyle f={\frac {1}{T}}}

, the frequency of a simple harmonic oscillator is

[
a
t
e
x
]
f
=

1
T

s
q
r
t

(

f
r
a
c
k
m

)

{\displaystyle f={\frac {1}{2\pi }}{\sqrt {\frac {k}{m}}}}\!\!

. Note that neither *T* nor *f* has any dependence on amplitude. Find two identical wooden or plastic rulers. Tape one end of each ruler firmly to the edge of a table so that the length of each ruler that protrudes from the table is the same. On the free end of one ruler tape a heavy object such as a few large coins. Pluck the ends of the rulers at the same time and observe which one undergoes more cycles in a time period, and measure the period of oscillation of each of the rulers. If the shock absorbers in a car go bad, then the car will oscillate at the least provocation, such as when going over bumps in the road and after stopping (See Figure 2). Calculate the frequency and period of these oscillations for such a car if the car's mass (including its load) is 900 kg and the force constant (*k*) of the suspension system is 6.53 × 104 N/m. Figure 2. The bouncing car makes a wavelike motion. If the restoring force in the suspension system can be described only by Hooke's law, then the wave is a sine function. (The wave is the trace produced by the headlight as the car moves to the right.) Strategy The frequency of the car's oscillations will be that of a simple harmonic oscillator as given in the equation

[
a
t
e
x
]
f
=

1
T

s
q
r
t

(

f
r
a
c
k
m

)

{\displaystyle f={\frac {1}{2\pi }}{\sqrt {\frac {k}{m}}}}\!\!

. The mass and the force constant are both given. Solution Enter the known values of *k* and *m*:

[
a
t
e
x
]
d
i
s
p
l
a
y
s
t
y
l
e
(
f
)
=

f
r
a
c
1
1
2
π

s
q
r
t

(

f
r
a
c
k
m

)

=

f
r
a
c
1
1
2
π

s
q
r
t

(

f
r
a
c
6.53
t
i
m
e
s

10

4

N

/

m

1900
t
e
x
t
{
k
g
}

)

{\displaystyle {\displaystyle (f)={\frac {1}{2\pi }}{\sqrt {\frac {k}{m}}}={\frac {1}{2\pi }}{\sqrt {\frac {6.53\times 10^{4}{\text{ N/m}}}{900{\text{ kg}}}}}\!\!

 Calculate the frequency:

[
a
t
e
x
]
f
r
a
c
1
1
2
π

s
q
r
t

(

f
r
a
c
6.53
t
i
m
e
s

10

4

N

/

m

1900
t
e
x
t
{
k
g
}

)

=
1.3656
/
t
e
x
t
(
s

)

^
(
-
1
)
a
p
p
r
o
x
1.36
/
t
e
x
t
(
s

)

^
(
-
1
)
=
1.36
t
e
x
t
{
H
z
}

{\displaystyle {\displaystyle (f)={\frac {1}{2\pi }}{\sqrt {\frac {k}{m}}}={\frac {1}{2\pi }}{\sqrt {\frac {6.53\times 10^{4}{\text{ N/m}}}{900{\text{ kg}}}}}\!\!

 Calculate the period, but it is simpler to use the relationship

[
a
t
e
x
]
T
=

1
f

{\displaystyle T={\frac {1}{f}}\!\!

 and substitute the value just found for *f*:

[
a
t
e
x
]
d
i
s
p
l
a
y
s
t
y
l
e
(
T
)
=

f
r
a
c
1
1
f

=

f
r
a
c
1
1
1.356
/
t
e
x
t
(
s

)

^
(
-
2
)

=
1.3656
/
t
e
x
t
(
s

)

^
(
-
1
)
a
p
p
r
o
x
1.36
/
t
e
x
t
(
s

)

^
(
-
1
)
=
1.36
t
e
x
t
{
H
z
}

{\displaystyle {\displaystyle T={\frac {1}{f}}={\frac {1}{1.356/{\text{ Hz}}}}=0.738{\text{ s}}}\!\!

 Discussion The values of *T* and *f* both seem about right for a bouncing car. You can observe these oscillations if you push down hard on the end of a car and let go. The Link between Simple Harmonic Motion and Waves Figure 3. The vertical position of an object bouncing on a spring is recorded on a strip of moving paper, leaving a sine wave. If a time-exposure photograph of the bouncing car were taken as it drove by, the headlight would make a wavelike streak, as shown in Figure 2. Similarly, Figure 3 shows an object bouncing on a spring as it leaves a wavelike "trace of its position on a moving strip of paper. Both waves are sine functions. All simple harmonic motion is intimately related to sine and cosine waves. The displacement as a function of time *t* in any simple harmonic motion—that is, one in which the net restoring force can be described by Hooke's law, is given by

[
a
t
e
x
]
x
(
t
)
=
X
c
o
s

f
r
a
c
2
π
(
t
)
T

{\displaystyle {\text{x(t)}=X\cos {\frac {2\pi (t)}{T}}}\!\!

, where *X* is amplitude. At *t* = 0, the initial position is *x*0 = *X*, and the displacement oscillates back and forth with a period *T*. (When *t* = *T*, we get *x* = *X* again because

cos
⁡
2
π
=
1
)
,

{\displaystyle \cos 2\pi =1),}

 Furthermore, from this expression for *x*, the velocity *v* as a function of time is given by

[
a
t
e
x
]
v
(
t
)
=
−
v
m
a
x

s
i
n

l
e
f
t
(

f
r
a
c
2
π
(
t
)
T

)

{\displaystyle {\text{v(t)}=-v_{\text{max}}\sin {\left({\frac {2\pi (t)}{T}}\right)}\!\!

, where

[
a
t
e
x
]
v
m
a
x

=

v
m
a
x

s
i
n

l
e
f
t
(

f
r
a
c
2
π
(
X
)
T

)

=
X
s
q
r
t

(

f
r
a
c
k
m

)

{\displaystyle {\text{v}_{\text{max}}}={\frac {v_{\text{max}}}{\sin {\left({\frac {2\pi (X)}{T}}\right)}}=X{\sqrt {\frac {k}{m}}}}\!\!

. The object has zero velocity at maximum displacement—for example, *v* = 0 when *t* = 0, and at that time *x* = *X*. The minus sign in the first equation for *v*(*t*) gives the correct direction for the velocity. Just after the start of the motion, for instance, the velocity is negative because the system is moving back toward the equilibrium point. Finally, we can get an expression for acceleration using Newton's second law. [Then we have *x*(*t*), *v*(*t*), and *a*(*t*), the quantities needed for kinematics and a description of simple harmonic motion.] According to Newton's second law, the acceleration is

[
a
t
e
x
]
a
=

f
r
a
c
F
m

=

f
r
a
c
(
−
k
x
)
m

{\displaystyle {\text{a}}={\frac {F}{m}}={\frac {\mathrm {-kx} }{m}}\!\!

. So, *a*(*t*) is also a cosine function:

[
a
t
e
x
]
a
(
t
)
=
−
k
X

c
o
s

f
r
a
c
2
π
(
t
)
T

{\displaystyle {\text{a(t)}=-{\frac {kX}{m}}\cos {\frac {2\pi (t)}{T}}}\!\!

. Hence, *a*(*t*) is directly proportional to and in the opposite direction to *a*(*t*). Figure 4 shows the simple harmonic motion of an object on a spring and presents graphs of *x*(*t*), *v*(*t*), and *a*(*t*) versus time. Figure 4. Graphs of *x* and *v* versus *t* for the motion of an object on a spring. The net force on the object can be described by Hooke's law, and so the object undergoes simple harmonic motion. Note that the initial position has the vertical displacement at its maximum value *X*; *v* is initially zero and then negative as the object moves down; and the initial acceleration is negative, back toward the equilibrium position and becomes zero at that point. The most important point here is that these equations are mathematically straightforward and are valid for all simple harmonic motion. They are very useful in visualizing waves associated with simple harmonic motion, including visualizing how waves add with one another. Suppose you pluck a banjo string. You hear a single note that starts out loud and slowly quiets over time. Describe what happens to the sound waves in terms of period, frequency and amplitude as the sound decreases in volume. Solution Frequency and period remain essentially unchanged. Only amplitude decreases as volume decreases. Part 2 A babysitter is pushing a child on a swing. At the point where the swing reaches *x*, where would the corresponding point on a wave of this motion be located? Solution *x* is the maximum deformation, which corresponds to the amplitude of the wave. The point on the wave would either be at the very top or the very bottom of the curve. A realistic mass and spring laboratory. Hang masses from springs and adjust the spring stiffness and damping. You can even slow time. Transport the lab to different planets. A chart shows the kinetic, potential, and thermal energy for each spring. Click to run the simulation. Selected Solutions Simple harmonic motion is oscillatory motion for a system that can be described only by Hooke's law. Such a system is also called a simple harmonic oscillator. Maximum displacement is the amplitude *X*. The period *T* and frequency *f* of a simple harmonic oscillator are given by

[
a
t
e
x
]
T
=

2
π

s
q
r
t

(

f
r
a
c
m
k

)

{\displaystyle T=2\pi {\sqrt {\frac {m}{k}}}}\!\!

 and

[
a
t
e
x
]
f
=

f
r
a
c
1
1
2
π

s
q
r
t

(

f
r
a
c
k
m

)

{\displaystyle f={\frac {1}{2\pi }}{\sqrt {\frac {k}{m}}}}\!\!

, where *m* is the mass of the system. Displacement in simple harmonic motion as a function of time is given by

[
a
t
e
x
]
x
(
t
)
=
X
c
o
s
(
ω
t
)

{\displaystyle x(t)=X\cos(\omega t)}

 or

[
a
t
e
x
]
x
(
t
)
=
X
c
o
s
(
ω
t
−
ϕ
)

{\displaystyle x(t)=X\cos(\omega t-\varphi)}

, where

ω
=

2
π
T

{\displaystyle \omega ={\frac {2\pi }{T}}}

 and

ϕ
=

2
π
f

{\displaystyle \phi ={\frac {2\pi }{f}}}

. The velocity is given by

[
a
t
e
x
]
v
(
t
)
=
−
v
m
a
x

s
i
n
(
ω
t
)

{\displaystyle v(t)=-v_{\text{max}}\sin(\omega t)}

 or

[
a
t
e
x
]
v
(
t
)
=
−
v
m
a
x

s
i
n
(
ω
t
−
ϕ
)

{\displaystyle v(t)=-v_{\text{max}}\sin(\omega t-\varphi)}

, where

[
a
t
e
x
]
v
m
a
x

=

v
m
a
x

s
i
n
(
ω
t
)

{\displaystyle v_{\text{max}}={\frac {v_{\text{max}}}{\sin(\omega t)}}=X\omega }

. The acceleration is found to be

[
a
t
e
x
]
a
(
t
)
=
−
k
X

c
o
s
(
ω
t
)

{\displaystyle a(t)=-{\frac {kX}{m}}\cos(\omega t)}

 or

[
a
t
e
x
]
a
(
t
)
=
−
k
X

c
o
s
(
ω
t
−
ϕ
)

{\displaystyle a(t)=-{\frac {kX}{m}}\cos(\omega t-\varphi)}

. Conceptual Questions What conditions must be met to produce simple harmonic motion? (a) If frequency is not constant for some oscillation, can the oscillation be simple harmonic motion? (b) Can you think of any examples of harmonic motion where the frequency may depend on the amplitude? Give an example of a simple harmonic oscillator, specifically noting how its frequency is independent of amplitude. Explain why you expect an object made of a stiff material to vibrate at a higher frequency than a similar object made of a spongy material. As you pass a freight truck with a trailer on a highway, you notice that its trailer is bouncing up and down slowly. Is it more likely that the trailer is heavily loaded or nearly empty? Explain your answer. Some people modify cars to be much closer to the ground than when manufactured. Should they install stiffer springs? Explain your answer. Problems & Exercises A type of cuckoo clock keeps time by having a mass bouncing on a spring, usually something cute like a cherub in a chair. What force constant is needed to produce a period of 0.500 s for a 0.0150-kg mass? If the spring constant of a simple harmonic oscillator is doubled, by what factor will the mass of the system need to change in order for the frequency of the motion to remain the same? A 0.500-kg mass suspended from a spring oscillates with a period of 1.50 s. How much mass must be added to the object to change the period to 2.00 s? By how much leeway (both percentage and mass) would you have in the selection of the mass of the object in the previous problem if you did not wish the new period to be greater than 2.01 s or less than 1.99 s? Suppose you attach the object with mass *m* to a vertical spring originally at rest, and let it bounce up and down. You release the object from rest at the spring's original rest length. (a) Show that the spring exerts an upward force of 2.00 mg on the object at its lowest point. (b) If the spring has a force constant of 10.0 N/m and a 0.25-kg-mass object is set in motion as described, find the amplitude of the oscillations. (c) Find the maximum velocity. A diver on a diving board is undergoing simple harmonic motion. Her mass is 55.0 kg and the period of her motion is 0.800 s. The next diver is a male whose period of simple harmonic oscillation is 1.05 s. What is his mass if the mass of the board is negligible? Suppose a diving board with no one on it bounces up and down in a simple harmonic motion with a frequency of 4.00 Hz. The board has an effective mass of 10.0 kg. What is the frequency of the simple harmonic motion of a 75.0-kg diver on the board? The device pictured in Figure 6 entertains infants while keeping them from wandering. The child bounces in a harness suspended from a door frame by a spring constant. Figure 6. This child's toy relies on springs to keep infants entertained. (credit: "Humboldthead, Flickr") (a) If the spring stretches 0.250 m while supporting an 8.0-kg child, what is its spring constant? (b) What is the time for one complete bounce of this child? (c) What is the child's maximum velocity if the amplitude of her bounce is 0.200 m? A 90.0-kg skydiver hanging from a parachute bounces up and down with a period of 1.50 s. What is the new period of oscillation when a second skydiver, whose mass is 60.0 kg, hangs from the legs of the first, as seen in Figure 7. Figure 7. The oscillations of one skydiver are about to be affected by a second skydiver. (credit: U.S. Army, www.army.mil) amplitude: the maximum displacement from the equilibrium position of an object oscillating around the equilibrium position simple harmonic motion: the oscillatory motion in a system where the net force can be described by Hooke's law simple harmonic oscillator: a device that implements Hooke's law, such as a mass that is attached to a spring, with the other end of the spring being connected to a rigid support such as a wall 1. 2.37 N/m 3. 0.389 kg 6. 94.7 kg 9. 1.94 s